

$$\frac{df}{d\eta} = Ae^{-\eta^2}$$

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$$f(\eta) = A \int_0^\eta e^{-\xi^2} d\xi + D$$

The B.C. $u(0, t) = U \Rightarrow UD = 1 \Rightarrow D = 1$

$$u(y, t) = U \left\{ 1 + A \int_0^y e^{-\xi^2} d\xi \right\} \quad \text{--- (3)}$$

The I.C. $u(y, 0) = 0$

$\forall y \Rightarrow f(\eta) \rightarrow 0$ when $\eta \rightarrow \infty$

$$0 = U \left\{ 1 + A \int_0^\infty e^{-\xi^2} d\xi \right\}$$

$$0 = U \left\{ 1 + A \frac{\sqrt{\pi}}{2} \right\}$$

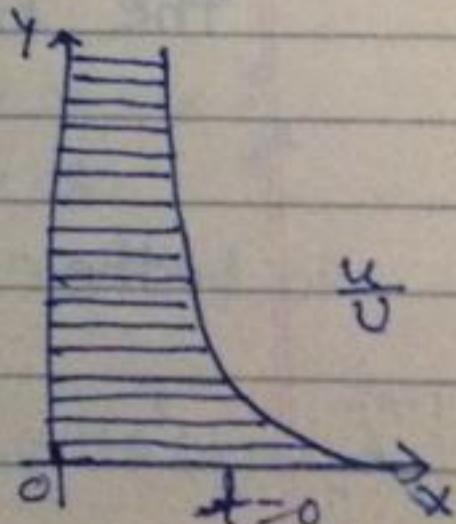
$$\Rightarrow A = -\frac{2}{\sqrt{\pi}} \quad \text{--- (4)}$$

equation (3) And (4) the velocity distribution

$$u(y, t) = U \left\{ 1 - \frac{2}{\sqrt{\pi}} \int_0^y e^{-\xi^2} d\xi \right\}$$

$$= U (1 - \operatorname{erf} \eta)$$

$$u(y, t) = U \operatorname{erfc} \eta$$



Ques:- Discuss the unsteady motion and dimensional velocity profile which are generated by similarity variable

$$= U\alpha^2 t^{-2n} f''$$

$$= U\alpha^2 t^{-2n} f''$$

$$\frac{\partial^2 y}{\partial t^2} = \frac{U\alpha^2}{t^{2n}} f'' \quad \rightarrow \textcircled{2}$$

From the relation \textcircled{2} And \textcircled{1} we have

$$-\frac{Un}{t} f' = v \frac{U\alpha^2}{t^{2n}} f''$$

The explicit time dependence may be eliminated.

By considering $n = \frac{1}{2}$ So we have

$$f'' + \frac{m}{v\alpha^2} f' = 0$$

$$\Rightarrow f'' + \left(\frac{n}{2v\alpha^2} \right) f' = 0 \quad \text{where } n = \frac{\alpha^2}{t^{1/2}}$$

The dimensions of $\left(\frac{y}{t^{1/2}}\right)$ are a length divided by the square root of time.
So α is taken, $\alpha = \frac{1}{\sqrt{v}}$

The similarity variable reduces

$$\eta = \frac{1}{2\sqrt{vt}}$$

The differential equation

$$f'' + 2\eta f' = 0$$

$$\frac{f''}{f'} = -2\eta$$

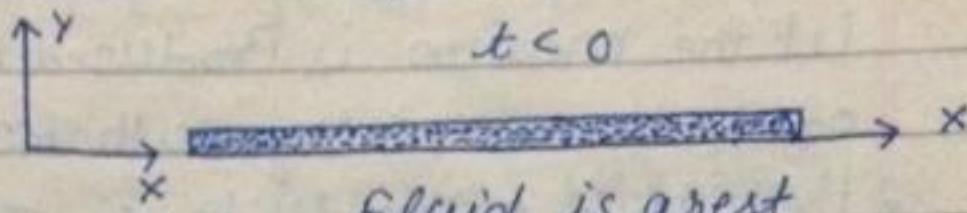
Integrating, we have

$$\log f' = -\eta^2 + \log A$$

Where $\log A$ is constant.

§ 11.7 Unsteady Motion of a flat Plate

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A semi infinite body of liquid with constant density ρ is bounded on one side by a flat surface. Initially the fluid and the solid surface are at rest. At time $t=0$ the direction x with the constant velocity U .

Equation of continuity $\frac{\partial u}{\partial x} = 0 \Rightarrow u = u(y, t)$
the velocity component $v = u(y, t)$, $\vartheta = 0$, $\omega = 0$,
there is no pressure gradient or gravity force in the X -direction.

the Navier - Stoke's equation.

$$\frac{\partial u}{\partial t} = v \frac{\partial^2 u}{\partial y^2} \quad \text{--- (1)}$$

I.C. $u(y, 0) = 0$ when $t, 0 \neq y$

B.C. $u(0, t) = U \quad \forall t > 0$

B.C. $u(\infty, t) = 0 \quad \forall t > 0$

Let $u(y, t) = U f(n)$

where $n = \alpha \frac{y}{t^n}$

the function $f(n, t)$ is known as similarity solution and α is a constant of proportionality.

$$\alpha \frac{\partial u}{\partial t} = \frac{\partial u}{\partial n} \cdot \frac{\partial n}{\partial t} = U s' (-n \alpha y t^{-n-1})$$

$$= \frac{Un}{t} n f'$$

$$\frac{\partial u}{\partial y^2} = U \alpha t^{-n} f'' \frac{\partial n}{\partial y}$$

• where $B = -\frac{A}{3}$

Let the motion is produced by a solid sphere of radius a rotating with angular velocity Ω

- the liquid extends to ∞ and is rest there
- We have:

$$\lambda = a, \omega = \Omega$$

$$\checkmark c = 0, B = a^3 \Omega$$

$$\boxed{\omega = \frac{a^3}{\lambda^3} \Omega} \quad \textcircled{11}$$

Again $\lambda = a, \omega = \Omega$

$$\lambda = b, \omega = 0$$

From $\textcircled{10}$ and $\textcircled{12}$ we have

$$\Omega = \frac{B}{a^3} + c \quad \text{and} \quad 0 = \frac{B}{b^3} + c$$

$$\Rightarrow B = \frac{a^3 b^3}{b^3 - a^3} \Omega, \quad c = -\left(\frac{a^3}{b^3 - a^3}\right) \Omega$$

Substituting the values of the constants in $\textcircled{10}$
we have

$$\boxed{\omega = \frac{a^3}{\lambda^3} \left(\frac{b^3 - \lambda^3}{b^3 - a^3} \right) \Omega} \quad \text{IV}$$

$= x = x$